

# Errata in “Statistical Mechanics in a Nutshell”

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## Chapter 2. Thermodynamics

### 2.17 Equations of state

- **Page 40, line 12 from bottom. Read:**  
obtained by deriving  
**correct to:**  
obtained by differentiating

## Chapter 3. The Fundamental Postulate

### 3.5 Quantum States

- **Page 66, line 2. Read:**  
moving in one dimension along a segment of length  $n$ .  
**correct to:**  
moving in one dimension along a line of length  $L$ .
- **Page 66, line 3. Read:**  
possible energy values are  $E_n = \hbar^2 \pi^2 n^2 / L^2$   
**correct to:**  
possible energy values are  $E_n = \hbar^2 \pi^2 n^2 / (2mL^2)$

### 3.18 Fluctuations of Uncorrelated Particles

- **Page 88. From the end of line 4 to the end of the section. Read:**  
In fact, one has...  
... Therefore,

$$p = \frac{k_B T}{v} = \frac{N k_B T}{V}. \quad (3.153)$$

**correct to:**  
In fact, one has

$$\langle N^2 \rangle - \langle N \rangle^2 = \left. \frac{\partial^2 \ln Z_{GC}}{\partial (\mu/k_B T)^2} \right)_{T,V} = k_B T \left. \frac{\partial N}{\partial \mu} \right)_{T,V}. \quad (3.150)$$

On the other hand, since  $\mu$  is an intensive variable, function of  $T$  and of the extensive variables  $V$  and  $N$ , one has the Euler equation

$$N \left( \frac{\partial \mu}{\partial N} \right)_{T,V} + V \left( \frac{\partial \mu}{\partial V} \right)_{T,N} = 0. \quad (3.151)$$

Thus from equations (3.149–150) we obtain

$$k_B T = N \left( \frac{\partial \mu}{\partial N} \right)_{T,V} = -V \left( \frac{\partial \mu}{\partial V} \right)_{T,N} = V \left( \frac{\partial p}{\partial N} \right)_{T,V}, \quad (3.152)$$

where we have exploited a Maxwell relation. Integrating this equation with respect to  $N$ , with the obvious boundary condition  $p(N=0) = 0$  yields

$$pV = N k_B T. \quad (3.153)$$

## Chapter 4. Interaction-Free systems

### 4.1 Harmonic Oscillators

#### 4.1.1 The Equipartition Theorem

- Page 90, line 16f. Read: positive definitive  
correct to: positive definite

### 4.3 Boson and Fermion Gases

#### 4.3.1 Electrons in Metals

- Page 109, line 5, eq. (4.96). Read:

$$C_V = \frac{\pi V^2}{3} k_B T \omega(\epsilon_F),$$

correct to:

$$C_V = \frac{\pi V^2}{3} k_B^2 T \omega(\epsilon_F),$$

### 4.4 Einstein condensation

- Page 114. Caption to figure 4.7, second line. Read: the rescaled density  $p\lambda^3$ .  
correct to: the rescaled density  $\rho\lambda^3$ .

#### 4.5.1 Myoglobin and Hemoglobin

- Page 116, line 6 from bottom. Read:  $\sum_{\alpha=1}^{N/4} \sum_{i=1}^4 \tau_{\alpha i}$  of adsorbed molecules  
correct to:  $\sum_{\alpha=1}^{N/4} \sum_{i=1}^4 \langle \tau_{\alpha i} \rangle$  of adsorbed molecules

## Chapter 6. Renormalization Group

### Relevant and Irrelevant Operators

- Page 184, line 10, eq. (6.57). Read:

$$\langle \phi_0 \phi_{\mathbf{r}} \rangle_{\mathcal{H}} = b^{2d} \zeta^{-2} \langle \phi'_0 \phi'_{\mathbf{r}/b} \rangle_{\mathcal{H}'}$$

correct to:

$$\langle \phi_0 \phi_{\mathbf{r}} \rangle_{\mathcal{H}} = b^{-2d} \zeta^{-2} \langle \phi'_0 \phi'_{\mathbf{r}/b} \rangle_{\mathcal{H}'}$$

- Page 184, line 10, eq. (6.57). Read:

correct to:

$$d + 2 - \eta = 2 \frac{\ln \zeta}{\ln b}$$

correct to:

$$d + 2 - \eta = -2 \frac{\ln \zeta}{\ln b}$$

## 6.6 Renormalization in Fourier Space

### 6.6.1 Introduction

- Page 190, line 12 from bottom, eq. (6.91). Read:

$$\phi_i = \sum_{\mathbf{k}} \phi_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{r}_i}$$

correct to:

$$\phi_i = \sum_{\mathbf{k}} \phi_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{r}_i}$$

- Page 190, line 6 from bottom. Read:  
For a simple cubic lattice, as we saw in chapter 2, one has  
correct to:  
For a simple cubic lattice, as we saw in chapter 5, one has

### 6.6.2 Gaussian Model

- Page 193, line 8. Read:  
coefficients of  $\kappa^n$  with  $n \neq 0$   
correct to:  
coefficients of  $k^n$  in  $\Delta(\mathbf{k})$  with  $n \neq 0$

### 6.6.4 Critical Exponents to First Order in $\epsilon$

- Page 199, line 5. Read:  
(as we shall from now on)  
correct to:  
(as we shall set from now on)

## Chapter 7. Classical Fluids

### 7.2 Reduced Densities

#### 7.2.3 Measure of $g(r)$

- Page 225, line 7. Read:  
define the factor structure  $S(\mathbf{k})$ :  
correct to:  
define the structure factor  $S(\mathbf{k})$ :

#### 7.2.4 BBGKY Hierarchy

- Page 225, line 7 from bottom, Eq. (7.47). Read:

$$\rho^{(2)}(\mathbf{x}_1, \mathbf{x}_2) = \rho^{(2)} g(\mathbf{x}_1, \mathbf{x}_2),$$

correct to:

$$\rho^{(2)}(\mathbf{x}_1, \mathbf{x}_2) = \rho^2 g(\mathbf{x}_1, \mathbf{x}_2),$$

### 7.3 Virial Expansion

- Page 229, line 16. Read:  
**Exercise 7.2** By comparing (7.70) with (7.12), show that in this approximation the  $g(r)$  is expressed by

$$g(r) = 1 + f(r).$$

correct to:

**Exercise 7.2** Show that in the present approximation, the  $g(r)$  is expressed by

$$g(r) = 1 + f(r),$$

where  $f(r)$  is the Mayer function, and that therefore equations (7.12) and (7.70) are compatible.

- Page 230, line 15. Read:  
3. Express this quantity as a function of the second virial coefficient  $B_2(T)$  and evaluate the inversion temperature  $T^*$  in which  $\partial T/\partial p)_H$  changes sign.  
correct to:  
3. By expressing  $\partial V/\partial T)_p$  as a function of the second virial coefficient  $B_2(T)$ , evaluate the inversion temperature  $T^*$  in which  $\partial T/\partial p)_H$  changes sign.

### 7.3.1 Higher Virial Coefficients

- Page 235, line 4, eq. (7.96). Read:

$$\frac{p}{\rho k_B T} = \frac{1 + \eta + \eta^2 - \eta^3}{(1 - \eta)^3}.$$

correct to:

$$\frac{p}{\rho k_B T} = \frac{1 + \eta + \eta^2 - \eta^3}{(1 - \eta)^3}.$$

### Convergence of the Fugacity and Virial Expansion

- Page 241, line 11, eq. (7.129). Read:

$$C(T) = \int d\mathbf{x} \left( e^{-u(\mathbf{x})/k_B T} - 1 \right) < +\infty.$$

correct to:

$$C(T) = \int d\mathbf{x} \left| e^{-u(\mathbf{x})/k_B T} - 1 \right| < +\infty.$$

## Chapter 8. Numerical Simulation

### 8.2 Molecular Dynamics

#### 8.2.2 Verlet Algorithm

- Page 258, line 9 from bottom. Read:  
and  $\tau = \sqrt{m/\epsilon_0}$  as a time scale,  
correct to:  
and  $\tau = \sqrt{mr_0^2/\epsilon_0}$  as a time scale,

### 8.4 Monte Carlo Method

#### 8.4.1 Markov Chains

- Page 264, line 7. Read:  
if, given any three different states  $a, b, c \in Q$ , one has

$$W_{ab}W_{bc}W_{ca} = W_{ac}W_{cb}W_{ba}. \quad (8.39)$$

correct to:

if, given any  $k$  different states  $i_1, i_2, \dots, i_k \in Q$  ( $k \geq 3$ ), one has

$$W_{i_1 i_2} W_{i_2 i_3} \cdots W_{i_k i_1} = W_{i_1 i_k} W_{i_k i_{k-1}} \cdots W_{i_2 i_1}. \quad (8.39)$$

## Chapter 9. Dynamics

### 9.9 Response Functions

- Page 298, line 12. Read:  
 $h(t) = h \delta(t - t')$ .  
correct to:  
 $h(t) = h \delta(t - t_0)$ .
- Page 298, line 13, eq. (9.111). Read:  
 $\langle X(t) \rangle = h \chi(t, t')$ .  
correct to:  
 $\langle X(t) \rangle = h \chi(t, t_0)$ .
- Page 299, line 8 from bottom. Read:  
 $x_{ij}(t)$ . For  $t > 0$ ,  
correct to:  
 $\chi_{ij}(t)$ . For  $t > 0$ ,

### 9.13 Variational Principle

- Page 306, line 1st from bottom. Read:  
an affinity  $F_i$   
correct to:  
an affinity  $F_1$
- Page 307, line 1. Read:  
Then, the result we just obtained that the stationary state...  
correct to:  
Then, the result we just obtained implies that the stationary state...
- Page 307, line 2. Read:  
In fact since...  
... and we obtain  $\partial \dot{S} / \partial X_k = 0$  for  $J_k = 0$  ( $k \neq 1$ ).  
correct to:  
In fact, upon a variation ( $\delta F_i$ ) of the forces, we obtain from equation (9.172) the corresponding variation of  $\dot{S}$

$$\delta \dot{S} = \sum_{ij} L_{ij} F_i \delta F_j = \sum_j J_j \delta F_j, \quad (9.179)$$

where we have used the relation (9.160). The term with  $j = 1$  vanishes because  $F_1$  is kept fixed. Thus  $\delta \dot{S} = 0$  implies  $J_k = 0$  for  $k \neq 1$ , since the  $\delta F_k$  with  $k \neq 1$  are arbitrary.

## Chapter 10. Complex systems

### 10.2. Percolation

#### 10.2.1 Analogy with Magnetic Phenomena

- Page 322. Eq. (10.47). Read:

$$\sum_s s \nu_s + P(p) = 1.$$

correct to:

$$\sum_s s\nu_s + P(p) = p.$$

- Page 322, line 3 from bottom. Read:

$$\chi = \frac{J}{k_{\text{B}}T} \sum_s s^2 \nu_s(p).$$

correct to:

$$\chi = \frac{1}{k_{\text{B}}T} \sum_s s^2 \nu_s(p).$$

### 10.2.1 Percolation in One Dimension

- Page 324. Eq. (10.54). Read:

$$S(p) = \frac{\sum_s s^2 \nu_s(p)}{\sum_s \nu_s(p)}.$$

correct to:

$$S(p) = \frac{\sum_s s \nu_s(p)}{\sum_s \nu_s(p)}.$$

### 10.2.3 Percolation on the Bethe lattice

- Page 326. Eq. (10.63). Read:

$$S(p) = p \frac{1 - (\zeta - 2)p}{1 - (\zeta - 1)p}, \quad \text{for } p < p_c,$$

correct to:

$$S(p) = p \frac{1 + p}{1 - (\zeta - 1)p}, \quad \text{for } p < p_c,$$

- Page 327, line 2. Read:

while  $\sum_s s\nu_s = 1$ ,

correct to:

while  $\sum_s s\nu_s = p$ ,

## 10.3. Disordered systems

### 10.3.3 Random Energy Model

- Page 344, line 13. Read:

with  $\epsilon \ll |E_c|$ .

correct to:

with  $|\epsilon| \ll |E_c|$ .

### 10.3.5 The replica method

- Page 349. Eq. (10.168). Read:

$$f = f_0 = k_B T \ln 2 - \left( \frac{J_0^2}{4k_B T} \right)$$

correct to:

$$f = f_0 = -k_B T \ln 2 - \left( \frac{J_0^2}{4k_B T} \right)$$

- Page 349, line 2 from bottom. Read:  
The minimum of this free energy is obtained when  
correct to:  
The extremum of this free energy is obtained when
- Page 350. Second line. Add the following sentence after eq. (10.173):  
One may check that the free energy per spin reaches a *maximum*, rather than a minimum, at this value of  $m$ . This is just one of the many surprises which appear in the replica method.

## Appendix

### B. Saddle Point Method

#### 0.0.1 B.1 Euler Integrals and the Saddle Point Method

- Page 366, line 3. Read:  
integral we studied before, slowly changing factors.  
correct to:  
integral we studied before, up to slowly changing factors.

### C. A Probability Refresher

#### C.2 Random Variables

- Page 370, line 9 from bottom. Read:

$$P(x) = \frac{1}{6} - \sum_{k=1}^6 \delta(x - k).$$

correct to:

$$P(x) = \frac{1}{6} \sum_{k=1}^6 \delta(x - k).$$

#### C.6 Central Limit Theorem

- Page 373, line 5, eq. (C.23), first line. Read:

$$\langle \exp(ik\bar{x}) \rangle = \left\langle \exp \left[ ik \frac{1}{N} \left( \sum_{i=1}^N \right) \right] \right\rangle = \left\langle \exp \left( \frac{ikx}{N} \right) \right\rangle$$

correct to:

$$\langle \exp(ik\bar{x}) \rangle = \left\langle \exp \left[ ik \frac{1}{N} \left( \sum_{i=1}^N x_i \right) \right] \right\rangle = \left\langle \exp \left( \frac{ikx}{N} \right) \right\rangle^N$$

#### D. Markov Chains

- **Page 337, line 4. Read:**  
Let  $\nu_k^{(\lambda)}$  be a right eigenvalue of W  
**correct to:**  
Let  $\nu_k^{(\lambda)}$  be a right eigenvector of W

#### E. Fundamental Physical Constants

- **Page 380, line 2. Read:**  
 $\hbar = h/(2\pi)$   
**correct to:**  
 $\hbar = h/(2\pi)$